

# Asymptotic expansions of oscillatory integrals with complex phase

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**ABSTRACT.** We consider saddle point integrals in  $d$  variables whose phase function is neither real nor purely imaginary. Results analogous to those for Laplace (real phase) and Fourier (imaginary phase) integrals hold whenever the phase function is analytic and nondegenerate. These results generalize what is well known for integrals of Laplace and Fourier type. The method is via contour shifting in complex  $d$ -space. This work is motivated by applications to asymptotic enumeration.

## 1. Introduction

Integrals of the form

$$(1.1) \quad \mathcal{I}(\lambda) := \mathcal{I}(\lambda; \phi, A) := \int e^{-\lambda\phi(x)} A(x) dx$$

arise in many areas of mathematics. There are many variations. This integral can arise in one or more variables; the variables may be real or complex; the integral may be global or taken in over a small neighborhood or oddly shaped set; varying degrees of smoothness may be assumed; varying degrees of degeneracy may be allowed near the critical points of the phase function,  $\phi$ . Often what is sought is a leading order estimate of  $\mathcal{I}(\lambda)$  as the positive real parameter  $\lambda$  tends to  $\infty$ , or an asymptotic series  $\mathcal{I}(\lambda) \sim \sum_n c_n g_n(\lambda)$  where  $\{g_n\}$  is a given sequence of elementary functions and the expansion is possibly nowhere convergent but satisfies

$$(1.2) \quad \mathcal{I}(\lambda) - \sum_{n=0}^{N-1} c_n g_n(\lambda) = O(g_N(\lambda))$$

for any  $N \geq 1$ .

In recent work on the asymptotics of multivariate generating functions [PW02; PW04; PW08; BBBP08; BP08; BP04], we have required results of this type in the case where the phase function  $\phi$  and the amplitude function  $A$  are analytic functions of several variables. The phase function is typically neither real nor

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purely imaginary, but satisfies  $\operatorname{Re}\{\phi\} \geq 0$ . Although these results are known, at least at a folklore level, we could not find this case analyzed in the literature in the generality we required. We therefore wrote this note to patch the gap and to remedy what would have been ghost citations in [PW04, Lemmas 4.7 and 4.8].

To better explain the relation between the present results and the literature, we discuss two special cases in which the results are well known, but by substantially different methods.

**Integrals of Fourier type.** Let  $f, A : \mathbb{R}^d \rightarrow \mathbb{R}$  be smooth (that is,  $C^\infty$ ) functions of  $d$  real variables, with  $A$  having compact support. Taking  $\phi = -if$  gives the Fourier-type integral

$$\mathcal{I}(\lambda) = \int e^{i\lambda f(x)} A(x) dx.$$

The standard method of studying this integral is as follows. If  $f$  has no critical points in the support of  $A$ , then integration by parts shows  $\mathcal{I}(\lambda)$  to be rapidly decreasing:  $\mathcal{I}(\lambda) = O(\lambda^{-N})$  for all positive integers,  $N$ . Using a partition of unity, the integral may therefore be localized to neighborhoods of the critical points of  $f$ . At an isolated critical point  $\nabla f$  vanishes; if the Hessian matrix is non-degenerate then the Morse lemma produces a smooth change of variables under which  $f(x) = S_\pm(x) := \sum_{j=1}^d \pm x_j^2$ . This reduces the general problem to the case  $\phi = iS$ . To solve this, expand  $A$  in a series. Each term may be explicitly integrated, resulting in an expansion in decreasing powers of  $\lambda$ :

$$(1.3) \quad \mathcal{I}(\lambda) = \sum_{n \geq 0} c_n \lambda^{-d/2-n}.$$

To see that the resulting series for  $\mathcal{I}(\lambda)$  satisfies (1.2), one uses integration by parts again to bound the remainder term. The coefficients  $\{c_n\}$  are determined by the derivatives of  $A$  and  $\phi$ . A particularly lucid treatment of this may be found in the first two sections of [Ste93, Chapter VIII].

**Integrals of Laplace type.** The other well studied case is the Laplace-type integral, where  $\phi$  is real. Localization of the integral to the minima of  $\phi$  is immediate because the integrand is exponentially small elsewhere. Integrating over balls whose radius has order  $\lambda^{-1/2}$  shows that near a quadratically nondegenerate minimum,  $x_0$ , the bound  $A(x) = O(|x - x_0|^N)$  translates into the bound

$$(1.4) \quad \mathcal{I}(\lambda) = O(\lambda^{-(d+N)/2}).$$

Again, one changes variables, expands  $A$  into a power series, and integrates term by term, to obtain the series

$$(1.5) \quad \mathcal{I}(\lambda) = \sum_{n \geq 0} c_n \lambda^{-(d+n)/2}.$$

Applying (1.4) to the integral of the remainder term shows that (1.5) is an asymptotic expansion for  $\mathcal{I}(\lambda)$ . Classical treatments of integrals of Laplace type may be found in many places [BH86; Won89], often accompanied by separate treatments of the Fourier case.

**Complex methods.** The series (1.3) and (1.5) are formally identical. Further inspection of formulae for  $c_n$  in the literature shows these to be nearly the same as well, differing only by constant factors of unit modulus. This points to the possibility of unifying the results and generalizing to arbitrary complex functions. An assumption of analyticity is required; technically, this is stronger than smoothness but in practice one is never satisfied without the other. Assuming analyticity, derivations in the Fourier and Laplace cases may indeed be unified via a hybrid approach. In one variable, this is carried out all the time under various names such as “steepest descent”, “stationary phase” or “saddle point”. Writing  $\mathcal{I}(\lambda)$  as a complex contour integral, the critical point will be a saddle point in  $\mathbb{C}^1$  for the real part of the phase function; the contour may then be re-oriented to pass through the saddle in the direction of steepest descent of  $-\phi$ , converting the integral into one of Fourier type and explaining why the series are nearly identical. This is carried out, for example, in [dB81] or [BH86, Chapter 7].

In more than one variable, complex phase results generalizing (1.3) and (1.5) are strangely absent from the literature. They are not treated in [Ste93; Hen91; BH86; Won89; Bre94; dB81]. We have found only two significant treatments of complex phases. One is Section 7.7 of [Hör83]. Here, notions from several complex variable theory are used to validate integration by parts, leading to asymptotic expansions in the case when  $A$  is smooth and compactly supported. This work has been used by [RW08] to produce explicit higher order terms in various expansions. As stated, Hörmander’s method does not allow integration over manifolds with boundary. The second treatment is [Fed89], which uses similar methodology to ours (complex contour deformation). The cited work omits the details, referring to the original article [Fed77] in Russian. Regions of integration with boundaries are allowed here, but only if  $\operatorname{Re}\{\phi\}$  is strictly maximized on the interior (our preliminary case, discussed in Section 4). In applications, work requiring high-dimensional saddle integration with complex phases is usually done by hand; see, e.g., [Wor04] or [BFSS00].

There are several possible reasons for this gap. One is that fewer people are aware of the corresponding techniques in several complex variables. Results such as Cauchy’s integral theorem in several variables and Stokes’ theorem for holomorphic  $d$ -forms in  $\mathbb{C}^d$ , while known to first-year graduate students in several complex variables, are not always known to those applying these techniques. The few mathematicians who do use these techniques to evaluate integrals [BH91; Lic91] are not unusually interested in anything as pedestrian as the generalization of standard saddle point integral evaluation. Secondly, even if one knows the techniques, moving contours in higher dimensions is a tricky business. The geometry can be much more complicated. When the real part of the phase vanishes at points that are not critical, further care must be taken in how the contour is moved. Finally, the need for such a generalization, illuminating though it may be, may have been too infrequent.

The organization of rest of the paper is as follows. In the next section we give some definitions having to do with stratified spaces. We then state our main result, Theorem 2.3. Section 3 records some easy computations in the case  $\phi$  is the standard phase function  $S(\mathbf{x}) := \sum_{j=1}^d x_j^2$ . Section 4 handles the general case under the assumption that the real part of  $\phi$  has a strict minimum at the critical point. Theorem 2.3 is proved in Section 5. Section 6 gives an application from [PW04] to

the estimation of coefficients of a bivariate generating function. Finally, Section 7 discusses further research directions.

## 2. Notation and statement of results

**Stratified spaces.** Because of the useful properties ensuing from the definition, we will use Whitney stratified spaces as our chains of integration. Aside from these useful properties, the details of the definition need not concern us, though for completeness we give a precise definition. Let  $I$  be a finite partially ordered set and define an  $I$ -decomposition of a topological space  $Z$  to be a partition of  $Z$  into a disjoint union of sets  $\{S_\alpha : \alpha \in I\}$  such that

$$S_\alpha \cap \overline{S_\beta} \neq \emptyset \iff S_\alpha \subseteq \overline{S_\beta} \iff \alpha \leq \beta.$$

**DEFINITION 2.1** (Whitney stratification). *Let  $Z$  be a closed subset of a smooth manifold  $\mathcal{M}$ . A **Whitney stratification** of  $Z$  is an  $I$ -decomposition such that*

- (i) *Each  $S_\alpha$  is a manifold in  $\mathbb{R}^n$ .*
- (ii) *If  $\alpha < \beta$ , if the sequences  $\{x_i \in S_\beta\}$  and  $\{y_i \in S_\alpha\}$  both converge to  $y \in S_\alpha$ , if the lines  $l_i = \overline{x_i y_i}$  converge to a line  $l$  and the tangent planes  $T_{x_i}(S_\beta)$  converge to a plane  $T$  of some dimension, then both  $l$  and  $T_y(S_\alpha)$  are contained in  $T$ .*

For example, any manifold is a Whitney stratified space with one stratum; any manifold with boundary is a Whitney stratified space with two strata, one being the interior and one the boundary; a  $k$ -simplex is a Whitney stratified space whose strata are all its faces. Whitney stratified spaces are closed under products and the set of products of strata will stratify the product. Every algebraic variety admits a Whitney stratification, although the singular locus filtration may be too coarse to be a Whitney stratification.

**Critical points.** Associated with the definition of a stratification is the stratified notion of a critical point. Observe that under this definition, any zero-dimensional stratum of  $\mathcal{M}$  is a critical point of  $\mathcal{M}$ .

**DEFINITION 2.2** (smooth functions and their critical points). *Say that a function  $\phi : \mathcal{M} \rightarrow \mathbb{C}$  on a stratified space  $\mathcal{M}$  is smooth if it is smooth when restricted to each stratum. A point  $p \in \mathcal{M}$  is said to be critical for the smooth function  $\phi$  if and only if the restriction  $d\phi|_S$  vanishes, where  $S$  is the stratum containing  $p$ .*

Let  $\mathcal{M} \subseteq \mathbb{C}^d$  be a real analytic,  $d$ -dimensional stratified space. This means that each stratum  $S$  is a subset of  $\mathbb{C}^d$  and each of the chart maps  $\psi$  from a neighborhood of the origin in  $\mathbb{R}^k$  to some  $k$ -dimensional stratum  $S \subseteq \mathbb{C}^d$  is analytic (the coordinate functions are convergent power series) with a nonsingular differential. It follows that  $\psi$  may be extended to a holomorphic map on a neighborhood of the origin in  $\mathbb{C}^k$ , whose range we denote by  $S \otimes \mathbb{C}$ . Choosing a small enough neighborhood, we may arrange for  $S \otimes \mathbb{C}$  to be a complex  $k$ -manifold embedded in  $\mathbb{C}^d$ . We say a function  $f : \mathcal{M} \rightarrow \mathbb{C}$  is analytic if it has a convergent power series expansion in a neighborhood of every point; an analytic function on  $\mathcal{M}$  may be extended to a complex analytic function on a neighborhood of  $\mathcal{M}$  in  $\mathcal{M} \otimes \mathbb{C} := \bigcup_{\alpha \in I} S_\alpha \otimes \mathbb{C}$ . Because we are interested in the integrals of  $d$ -forms over  $\mathcal{M}$ , there is no loss of generality in assuming that  $\mathcal{M}$  is contained in the closure of its  $d$ -dimensional strata, whence  $\mathcal{M} \otimes \mathbb{C}$  is a neighborhood of  $\mathcal{M}$  in  $\mathbb{C}^d$ . Real  $d$ -manifolds in  $\mathbb{C}^d$  are not naturally oriented, so we must assume that an orientation is given for

each  $d$ -stratum of  $\mathcal{M}$ , meaning that the chart maps from  $\mathbb{R}^d$  to  $\mathcal{M}$  must preserve orientations. The critical point  $p \in \mathcal{M}$  is said to be **quadratically nondegenerate** if  $p$  is in a  $d$ -dimensional stratum and  $\mathcal{H}(p)$  is nonsingular, where

$$\mathcal{H}(p) := \left( \frac{\partial^2 \phi}{\partial x_i \partial x_j} \right)_{1 \leq i, j \leq d}$$

is the Hessian matrix for the function  $\phi$  of Definition 2.2 on a neighborhood of  $p$  in  $\mathbb{C}^d$ .

Finally, we define a point of  $\mathcal{M}$  to be **stationary** if it is critical and  $\operatorname{Re}\{\phi\}$  vanishes there.

**Results.** The main result of the paper is an asymptotic expansion for  $\mathcal{I}(\lambda)$ .

**THEOREM 2.3.** *Let  $\mathcal{M}$  be a compact, real analytic stratified space of dimension  $d$  embedded in  $\mathbb{C}^d$ . Let  $A$  and  $\phi$  be analytic functions on a neighborhood of  $\mathcal{M}$  and suppose  $\operatorname{Re}\{\phi\} \geq 0$  on  $\mathcal{M}$ . Let  $G$  be the subset of stationary points of  $\phi$  on  $\mathcal{M}$  and assume that  $G$  is finite and each stationary point is quadratically nondegenerate. Then the integral  $\mathcal{I}(\lambda) := \int_{\mathcal{M}} e^{-\lambda \phi(x)} A(x) dx$  has an asymptotic expansion*

$$\mathcal{I}(\lambda) \sim \sum_{\ell=0}^{\infty} c_{\ell} \lambda^{-(d+\ell)/2}.$$

If  $A$  is nonzero at some point of  $G$  then the leading term is given by

$$(2.1) \quad c_0 = (2\pi)^{-d/2} \sum_{x \in G} A(x) \omega(x, \lambda) (\det \mathcal{H}(x))^{-1/2}.$$

where  $\omega(x, \lambda)$  is the unit complex number  $e^{-\lambda \phi(x)}$ . The choice of sign of the square root on the right-hand side of (2.1) is determined by choosing any analytic chart map  $\psi$  for a neighborhood of  $x$  and defining

$$(\det \mathcal{H}(x))^{-1/2} := (\det \mathcal{H}(\phi \circ \psi))^{-1/2} \det J_{\psi}$$

where  $J_{\psi}$  is the Jacobian matrix for  $\psi$  at  $x$  and the  $-1/2$  power of the determinant on the right is the product of the inverses of principal square roots of the eigenvalues.

**COROLLARY 2.4.** *Assume the hypotheses of Theorem 2.3 except that some of the points of  $G$  may be in strata of dimension  $d-1$  with a neighborhood in  $\mathcal{M}$  diffeomorphic to a halfspace in  $\mathbb{R}^d$ . Then the same conclusion holds except that the summand in (2.1) corresponding to such a point  $x$  must be multiplied by  $1/2$ .*

### 3. Preliminary results for the standard phase function

For  $\mathbf{x} \in \mathbb{R}^d$ , let  $S(\mathbf{x}) := \sum_{j=1}^d x_j^2$  denote the standard quadratic form. We begin with a couple of results on integrals of Laplace type where the phase function is the standard quadratic and  $A$  is a monomial. In one dimension,

**PROPOSITION 3.1.**

$$\int_{-\infty}^{\infty} x^n e^{-\lambda x^2} dx = \frac{n!}{(n/2)! 2^n \sqrt{2\pi}} \lambda^{-1/2-n/2}$$

if  $n$  is even, while the integral is zero if  $n$  is odd.

PROOF. For odd  $n$  the result is obvious from the fact that the integrand is an odd function. For  $n = 0$  the result is just the standard Gaussian integral. By induction, assume now the result for  $n - 2$ . Integrating by parts to get the second line, we have

$$\begin{aligned} \int x^n e^{-\lambda x^2} dx &= \int \frac{-x^{n-1}}{2\lambda} (-2\lambda x e^{-x^2} dx) \\ &= \frac{n-1}{2\lambda} \int x^{n-2} e^{-x^2} dx \\ &= \frac{n-1}{2\lambda} \frac{1}{\sqrt{2\pi}} \frac{(n-2)!}{(n/2-1)! 2^{n-1}} \lambda^{1/2-n/2} \end{aligned}$$

by the induction hypothesis. This is equal to  $\lambda^{-1/2-n/2} (2\pi)^{-1/2} n! / ((n/2)! 2^n)$ , completing the induction.  $\square$

COROLLARY 3.2 (monomial integral). *Let  $\mathbf{r}$  be any  $d$ -vector of nonnegative integers and let  $\mathbf{x}^{\mathbf{r}}$  denote the monomial  $x_1^{r_1} \cdots x_d^{r_d}$ . Then*

$$\int_{\mathbb{R}^d} \mathbf{x}^{\mathbf{r}} e^{-\lambda S(\mathbf{x})} d\mathbf{x} = \beta_{\mathbf{r}} \lambda^{-(d+|\mathbf{r}|)/2}$$

where

$$(3.1) \quad \beta_{\mathbf{r}} := (2\pi)^{-d/2} \prod_{j=1}^d \frac{r_j!}{(r_j/2)! 2^{r_j}}$$

if all the components  $r_j$  are even, and zero otherwise.

PROOF. The integral factors into

$$\prod_{j=1}^d \left[ \int_{-\infty}^{\infty} x_j^{r_j} e^{-\lambda r_j^2} dx_j \right].$$

$\square$

To integrate term by term over a Taylor series for  $A$ , we need the following estimate.

LEMMA 3.3 (big-O estimate). *Let  $A$  be any smooth function satisfying  $A(\mathbf{x}) = O(|\mathbf{x}|^r)$  at the origin. Then the integral of  $A(\mathbf{x})e^{-\lambda S(\mathbf{x})}$  over any compact set  $K$  may be bounded from above as*

$$\int_K A(\mathbf{x}) e^{-\lambda S(\mathbf{x})} d\mathbf{x} = O(\lambda^{-(d+r)/2})$$

PROOF. Because  $K$  is compact and  $A(\mathbf{x}) = O(|\mathbf{x}|^r)$  at the origin, it follows that there is some constant  $c$  for which  $|A(\mathbf{x})| \leq c|\mathbf{x}|^r$  on all of  $K$ . Let  $K_0$  denote the intersection of  $K$  with the ball  $|\mathbf{x}| \leq \lambda^{-1/2}$  and for  $n \geq 1$  let  $K_n$  denote the intersection of  $K$  with the shell  $2^{n-1}\lambda^{-1/2} \leq |\mathbf{x}| \leq 2^n\lambda^{-1/2}$ . On  $K_0$  we have

$$|A(\mathbf{x})| \leq c\lambda^{-r/2}$$

while trivially

$$\int_{K_0} e^{-\lambda S(\mathbf{x})} d\mathbf{x} \leq \int_{K_0} d\mathbf{x} \leq c_d \lambda^{-d/2}.$$

Thus

$$\left| \int_{K_0} A(\mathbf{x}) e^{-\lambda S(\mathbf{x})} d\mathbf{x} \right| \leq c' \lambda^{-(r+d)/2}.$$

For  $n \geq 1$ , on  $K \cap K_n$ , we have the upper bounds

$$\begin{aligned} |A(\mathbf{x})| &\leq 2^{rn} c \lambda^{-r/2} \\ e^{-\lambda S(\mathbf{x})} &\leq e^{-2^{n-1}} \\ \int_{K_n} d\mathbf{x} &\leq 2^{rn} c_d \lambda^{-d/2}. \end{aligned}$$

Letting  $c'' := c \cdot c_d \cdot \sum_{n=1}^{\infty} 2^{2rn} e^{-2^{n-1}} < \infty$ , we may sum to find that

$$\sum_{n=0}^{\infty} \left| \int_{K_n} A(\mathbf{x}) e^{-\lambda S(\mathbf{x})} d\mathbf{x} \right| \leq (c' + c'') \lambda^{-(r+d)/2},$$

proving the lemma.  $\square$

It is now easy to compute a series for  $\mathcal{I}(\lambda)$  in the case where  $\phi$  is the standard quadratic and the integral is over a neighborhood of the origin in  $\mathbb{R}^d$ .

**THEOREM 3.4** (standard phase). *Let  $A(\mathbf{x}) = \sum_{\mathbf{r}} a_{\mathbf{r}} \mathbf{x}^{\mathbf{r}}$  be any analytic function defined on a neighborhood  $\mathcal{N}$  of the origin in  $\mathbb{R}^d$ . Let*

$$(3.2) \quad \mathcal{I}(\lambda) := \int_{\mathcal{N}} A(\mathbf{x}) e^{-\lambda S(\mathbf{x})} d\mathbf{x}.$$

Then

$$\mathcal{I}(\lambda) \sim \sum_n \left[ \sum_{|\mathbf{r}|=n} a_{\mathbf{r}} \beta_{\mathbf{r}} \right] \lambda^{-(n+d)/2}$$

as an asymptotic series expansion in decreasing powers of  $\lambda$ , with  $\beta_{\mathbf{r}}$  as in (3.1).

**PROOF.** Write  $A(\mathbf{x})$  as a power series up to degree  $N$  plus a remainder term:

$$A(\mathbf{x}) = \left( \sum_{n=0}^N \sum_{|\mathbf{r}|=n} a_{\mathbf{r}} \mathbf{x}^{\mathbf{r}} \right) + R(\mathbf{x})$$

where  $R(\mathbf{x}) = O(|\mathbf{x}|^{N+1})$ . Using Corollary 3.2 to integrate all the monomial terms and Lemma 3.3 to bound the integral of  $R(\mathbf{x}) e^{-\lambda S(\mathbf{x})}$  shows that

$$\mathcal{I}(\lambda) = \sum_{n=0}^N \sum_{|\mathbf{r}|=n} a_{\mathbf{r}} \beta_{\mathbf{r}} \lambda^{-(d+n)/2} + O(\lambda^{-(d+N)/2-1})$$

which proves the asymptotic expansion.  $\square$

#### 4. The case of a strict minimum

In this section, we continue to integrate over a neighborhood of the origin in  $\mathbb{R}^d$ , but we generalize to any analytic phase function  $\phi$  with the restriction that the real part of  $\phi$  have a strict minimum at the origin. The assumption of a strict minimum localizes the integral to the origin, so the only tricky aspects are keeping track of the sign (Lemma 4.3) and being rigorous about moving the contour.

**THEOREM 4.1.** *Let  $A$  and  $\phi$  be complex-valued analytic functions on a compact neighborhood  $\mathcal{N}$  of the origin in  $\mathbb{R}^d$  and suppose that the real part of  $\phi$  is nonnegative, vanishing only at the origin. Suppose that the Hessian matrix  $\mathcal{H}$  of  $\phi$  at the origin is nonsingular. Denoting  $\mathcal{I}(\lambda) := \int_{\mathcal{N}} A(\mathbf{x}) e^{-\lambda \phi(\mathbf{x})} d\mathbf{x}$ , there is an asymptotic expansion*

$$\mathcal{I}(\lambda) \sim \sum_{\ell \geq 0} c_\ell \lambda^{d/2 - \ell}$$

where

$$c_0 = A(\mathbf{0}) \frac{(2\pi)^{-d/2}}{\sqrt{\det \mathcal{H}}}$$

and the choice of sign is defined by taking the product of the principal square roots of the eigenvalues of  $\mathcal{H}$ .

The proof is essentially a reduction to the case of standard phase. The key is the well known Morse Lemma. The proof given in [Ste93, VIII:2.3.2] is for the smooth category and purely real or imaginary phase but extends without significant change to complex values and the analytic category. For completeness, we include the adapted proof.

**LEMMA 4.2 (Morse lemma).** *There is a bi-holomorphic change of variables  $\mathbf{x} = \psi(\mathbf{y})$  such that  $\phi(\psi(\mathbf{y})) = S(\mathbf{y}) := \sum_{j=1}^d y_j^2$ . The differential  $J_\psi = d\psi(\mathbf{0})$  will satisfy  $(\det J_\psi)^2 = (\det \frac{1}{2} \mathcal{H})^{-1}$ .*

**PROOF.** Addressing the second conclusion first, we recall how the Hessian matrix behaves under a change of variables. If  $\psi : \mathbb{C}^d \rightarrow \mathbb{C}^d$  is bi-holomorphic on a neighborhood of  $\mathbf{x}$  and if  $\phi$  has vanishing gradient at  $\psi(\mathbf{x})$  and Hessian matrix  $\mathcal{H}$  there, then the Hessian matrix  $\tilde{\mathcal{H}}$  of  $\phi \circ \psi$  at  $\mathbf{x}$  is given by

$$\tilde{\mathcal{H}} = J_\psi^T \mathcal{H} J_\psi$$

where  $J_\psi$  is the Jacobian matrix for  $\psi$  at  $\mathbf{x}$ . The standard form  $S$  has Hessian matrix equal to twice the identity, hence any function  $\psi$  satisfying  $\phi \circ \psi = S$  must satisfy

$$2 Id = J_\psi^T \mathcal{H} J_\psi.$$

Dividing by two and taking determinants yields  $|\mathbf{J}_\psi|^2 \det(\frac{1}{2} \mathcal{H}) = 1$ , proving the second conclusion.

To prove the change of variables, the first step is to write

$$\phi(\mathbf{x}) = \sum_{j,k=1}^d x_j x_k \phi_{j,k}$$

where the functions  $\phi_{j,k} = \phi_{k,j}$  are analytic and satisfy  $\phi_{j,k}(\mathbf{0}) = (1/2) \mathcal{H}_{j,k}$ . It is obvious from a formal power series viewpoint that this can be done because the summand  $x_j x_k \phi_{j,k}$  can be any power series with coefficients indexed by the orthant  $\{\mathbf{r} : \mathbf{r} \geq \delta_j + \delta_k\}$ ; these orthants cover  $\{\mathbf{r} : |\mathbf{r}| \geq 2\}$ , so we may obtain any function  $\phi$  vanishing to order two; matching coefficients on the terms of order precisely two shows that  $\phi_{j,k}(\mathbf{0}) = (1/2) \mathcal{H}_{j,k}$ .

More constructively, we may give a formula for  $\phi_{j,k}$ . There is plenty of freedom, but a convenient choice is to let  $a_{\mathbf{r}}$  denote the coefficient of  $\mathbf{x}^{\mathbf{r}}$  in  $\phi(\mathbf{x})$  and to take

$$x_k x_j \phi_{j,k}(\mathbf{x}) := \sum_{|\mathbf{r}| \geq 2} \frac{r_j(r_k - \delta_{j,k})}{|\mathbf{r}|(|\mathbf{r}| - 1)} a_{\mathbf{r}} \mathbf{x}^{\mathbf{r}}.$$



For fixed  $\mathbf{r}$ , it is easy to check that

$$\sum_{1 \leq j, k \leq d} \frac{r_j(r_k - \delta_{j,k})}{|\mathbf{r}|(|\mathbf{r}| - 1)} = 1$$

whence  $\phi = \sum x_j x_k \phi_{j,k}$ . Alternatively, the following analytic computation from [Ste93] verifies that  $\phi = \sum_{j,k} x_j x_k \phi_{j,k}$ . Any function  $f$  vanishing at zero satisfies  $f(t) = \int_0^1 (1-s)f'(s)ds$ , as may be seen by integrating by parts (take  $g(s) = -(1-s)$ ). Fix  $\mathbf{x}$  and apply this with  $f(t) = (d/dt)\phi(t\mathbf{x})$  to obtain

$$\phi(\mathbf{x}) = \int_0^1 \frac{d}{dt} \phi(t\mathbf{x}) dt = \int_0^1 (1-t) \frac{d^2}{dt^2} \phi(t\mathbf{x}) dt.$$

The multivariate chain rule gives

$$\frac{d^2}{dt^2} \phi(t\mathbf{x}) = \sum_{j,k} x_j x_k \frac{\partial^2 \phi}{\partial x_j \partial x_k}(t\mathbf{x});$$

plug in  $\phi = \sum_{\mathbf{r}} a_{\mathbf{r}} \mathbf{x}^{\mathbf{r}}$  and integrate term by term using  $\int_0^1 (1-t)t^{n-2} dt = \frac{1}{n(n-1)}$  to see that  $\phi = \sum_{j,k} x_j x_k \phi_{j,k}$ .

The second step is an induction. Suppose first that  $\phi_{j,j}(\mathbf{0}) \neq 0$  for all  $j$ . The function  $\phi_{1,1}^{-1}$  and a branch of the function  $\phi_{1,1}^{1/2}$  are analytic in a neighborhood of the origin. Set

$$y_1 := \phi_{1,1}^{1/2} \left[ x_1 + \sum_{k>1} \frac{y_k \phi_{1,k}}{\phi_{1,1}} \right].$$

Expanding, we find that the terms of  $y_1^2$  of total degree at most one in the terms  $x_2, \dots, x_d$  match those of  $\phi$  and therefore,

$$(4.1) \quad \phi(\mathbf{x}) = y_1^2 + \sum_{j,k \geq 2} x_j x_k h_{j,k}$$

for some analytic functions  $h_{j,k}$  satisfying  $h_{j,k}(\mathbf{0}) = (1/2) \mathcal{H}_{j,k}$ . Similarly, if

$$\phi(\mathbf{x}) = \sum_{j=1}^{r-1} y_j^2 + \sum_{j,k \geq r} x_j x_k h_{j,k}$$

then setting

$$y_r := \phi_{r,r}^{1/2} \left[ x_r + \sum_{k>r} \frac{y_k \phi_{r,k}}{\phi_{r,r}} \right]$$

gives

$$\phi(\mathbf{x}) = \sum_{j=1}^r y_j^2 + \sum_{j,k \geq r+1} x_j x_k \tilde{h}_{j,k}$$

for some analytic functions  $\tilde{h}_{j,k}$  still satisfying  $h_{j,k}(\mathbf{0}) = (1/2) \mathcal{H}_{j,k}$ . By induction, we arrive at  $\phi(\mathbf{x}) = \sum_{j=1}^d y_j^2$ , finishing the proof of the Morse Lemma in the case where each  $\mathcal{H}_{j,j}$  is nonzero.

Finally, if some  $\mathcal{H}_{j,k} = 0$ , because  $\mathcal{H}$  is nonsingular we may always find some unitary map  $U$  such that the Hessian  $U^T \mathcal{H} U$  of  $\phi \circ U$  has no vanishing diagonal entries. We know there is a  $\psi_0$  such that  $(\phi \circ U) \circ \psi_0 = S$ , and taking  $\psi = U \circ \psi_0$  finishes the proof in this case.  $\square$

PROOF OF THEOREM 4.1: The power series allows us to extend  $\phi$  to a neighborhood of the origin in  $\mathbb{C}^d$ . Under the change of variables  $\psi$  from the previous lemma, we see that

$$\begin{aligned}\mathcal{I}(\lambda) &= \int_{\psi^{-1}\mathcal{C}} A \circ \psi(\mathbf{y}) e^{-\lambda S(\mathbf{y})} (\det d\psi(\mathbf{y})) d\mathbf{y} \\ &:= \int_{\psi^{-1}\mathcal{C}} \tilde{A}(\mathbf{y}) e^{-\lambda S(\mathbf{y})} d\mathbf{y}\end{aligned}$$

for some analytic function  $\tilde{A}$ , where  $\mathcal{C}$  is a neighborhood of the origin in  $\mathbb{R}^n$ . We need to check that we can move the chain  $\psi^{-1}\mathcal{C}$  of integration back to the real plane.

Let  $h(\mathbf{z}) := \operatorname{Re}\{S(\mathbf{z})\}$ . The chain  $\mathcal{C}' := \psi^{-1}(\mathcal{C})$  lies in the region  $\{\mathbf{z} \in \mathbb{C}^d : h(\mathbf{z}) > 0\}$  except when  $\mathbf{z} = 0$ , and in particular,  $h \geq \epsilon > 0$  on  $\partial\mathcal{C}'$ . Let

$$H(\mathbf{z}, t) := \operatorname{Re}\{\mathbf{z}\} + (1 - t)i \operatorname{Im}\{\mathbf{z}\}.$$

In other words,  $H$  is a homotopy from the identity map to the map  $\pi$  projecting out the imaginary part of the vector  $\mathbf{z}$ . For any chain  $\sigma$ , the homotopy  $H$  induces a chain homotopy,  $H(\sigma)$  supported on the image of the support of  $\sigma$  under the homotopy  $H$  and satisfying

$$\partial H(\sigma) = \sigma - \pi\sigma + H(\partial\sigma).$$

With  $\sigma = \mathcal{C}'$ , observing that  $S(H(\mathbf{z}, t)) \geq S(\mathbf{z})$ , we see there is a  $(d + 1)$ -chain  $\mathcal{D}$  with

$$\partial\mathcal{D} = \mathcal{C}' - \pi\mathcal{C}' + \mathcal{C}''$$

and  $\mathcal{C}''$  supported on  $\{h > \epsilon\}$ . Stokes Theorem tells us that for any holomorphic  $d$ -form  $\omega$ ,

$$\int_{\partial\mathcal{D}} \omega = \int_{\mathcal{D}} d\omega = 0$$

and consequently, that

$$\int_{\mathcal{C}'} \omega = \int_{\pi\mathcal{C}'} \omega + \int_{\mathcal{C}''} \omega.$$

When  $\omega = \tilde{A}e^{-\lambda S} d\mathbf{y}$ , the integral over  $\mathcal{C}''$  is  $O(e^{-\lambda\epsilon})$ , giving

$$\mathcal{I}(\lambda) = \int_{\pi\mathcal{C}'} \tilde{A}(\mathbf{y}) e^{-\lambda S(\mathbf{y})} d\mathbf{y} + O(e^{-\lambda\epsilon}).$$

Up to sign, the chain  $\pi\mathcal{C}''$  is a disk in  $\mathbb{R}^d$  with the standard orientation plus something supported in  $\{h > \epsilon\}$ . To see this, note that  $\pi$  maps any real  $d$ -manifold in  $\mathbb{C}^d$  diffeomorphically to  $\mathbb{R}^d$  wherever the tangent space is transverse to the imaginary subspace. The tangent space to the support of  $\mathcal{C}'$  at the origin is transverse to the imaginary subspace because  $S \geq 0$  on  $\mathcal{C}'$ , whereas the imaginary subspace is precisely the negative  $d$ -space of the index- $d$  form  $S$ . The tangent space varies continuously, so in a neighborhood of the origin,  $\pi$  is a diffeomorphism. Observing that  $\tilde{A}(\mathbf{0}) = A(\mathbf{0}) \det(d\psi(\mathbf{0})) = A(\mathbf{0})(\det \mathcal{H})^{-1/2}$  and using Theorem 3.4 finishes the proof up to the choice of sign of the square root.

The map  $d\pi \circ d\psi^{-1}(\mathbf{0})$  maps the standard basis of  $\mathbb{R}^d$  to another basis for  $\mathbb{R}^d$ . Verifying the sign choice is equivalent to showing that this second basis is positively oriented if and only if  $\det(d\psi(\mathbf{0}))$  is the product of the principal square roots of the eigenvalues of  $\mathcal{H}$  (it must be either this or its negative). Thus we will be finished by applying the following lemma (with  $\alpha = \psi^{-1}$ ).

LEMMA 4.3. *Let  $W \subseteq \mathbb{C}^d$  be the set  $\{\mathbf{z} : \operatorname{Re}\{S(\mathbf{z})\} > 0\}$ . Pick any  $\alpha \in GL_d(\mathbb{C})$  mapping  $\mathbb{R}^d$  into  $\overline{W}$  and let  $M := \alpha^T \alpha$  be the matrix representing  $S \circ \alpha$ . Let  $\pi : \mathbb{C}^d \rightarrow \mathbb{R}^d$  be projection onto the real part. Then  $\pi \circ \alpha$  is orientation preserving on  $\mathbb{R}^d$  if and only if  $\det \alpha$  is the product of the principal square roots of the eigenvalues of  $M$  (rather than the negative of this).*

PROOF. First suppose  $\alpha \in GL_d(\mathbb{R})$ . Then  $M$  has positive eigenvalues, so the product of their principal square roots is positive. The map  $\pi$  is the identity on  $\mathbb{R}^d$  so the statement boils down to saying that  $\alpha$  preserves orientation if and only if it has positive determinant, which is true by definition. In the general case, let  $\alpha_t := \pi_t \circ \alpha$ , where  $\pi_t(\mathbf{z}) = \operatorname{Re}\{\mathbf{z}\} + (1-t)\operatorname{Im}\{\mathbf{z}\}$ . As we saw in the previous proof,  $\pi_t(\mathbb{R}^d) \subseteq \overline{W}$  for all  $0 \leq t \leq 1$ , whence  $M_t := \alpha_t^T \alpha_t$  has eigenvalues with nonnegative real parts. The product of the principal square roots of the eigenvalues is a continuous function on the set of nonsingular matrices with no negative real eigenvalues. The determinant of  $\alpha_t$  is a continuous function of  $t$ , and we have seen it agrees with the product of principal square roots of eigenvalues of  $M_t$  when  $t = 1$  (the real case), so by continuity, this is the correct sign choice for all  $0 \leq t \leq 1$ ; taking  $t = 0$  proves the lemma.  $\square$

For later use, we record one easy corollary of Theorem 4.1.

COROLLARY 4.4. *Assume the hypotheses of Theorem 4.1 and let  $\mathcal{N}'$  be the intersection of  $\mathcal{N}$  with a region diffeomorphic to a halfspace through the origin. If  $A(\mathbf{0}) \neq 0$  then*

$$\mathcal{I}'(\lambda) := \int_{\mathcal{N}'} A(\mathbf{x}) e^{-\lambda \phi(\mathbf{x})} d\mathbf{x} \sim \frac{c_0}{2} \lambda^{-d/2}$$

where  $c_0$  is the same as in the conclusion of Theorem 4.1.

PROOF. Under the change of variables  $\psi$  and the projection  $\pi$ , this region maps to a region  $\mathcal{N}''$  diffeomorphic to a halfspace with the origin on the boundary. Changing variables by  $\mathbf{y} = \lambda^{-1/2} \mathbf{x}$  and writing  $\mathcal{N}_\lambda$  for  $\lambda^{1/2} \mathcal{N}''$ , we have

$$\mathcal{I}'(\lambda) = \lambda^{-d/2} \int_{\mathcal{N}_\lambda} A_\lambda(\mathbf{y}) e^{-S(\mathbf{y})} d\mathbf{y}$$

where  $A_\lambda(\mathbf{y}) = (A \circ \psi)(\lambda^{-1/2} \mathbf{y})$ . The function  $A_\lambda$  converges to  $A(\mathbf{0})$  pointwise but also in  $L^2(\mu)$  where  $\mu$  is the Gaussian measure  $e^{-S(\mathbf{x})} d\mathbf{x}$ . Also, the regions  $\mathcal{N}_\lambda$  converge to a halfspace  $H$  in the sense that their indicators  $\mathbf{1}_{\mathcal{N}_\lambda}$  converge to  $\mathbf{1}_H$  in  $L^2(\mu)$ . Thus  $A_\lambda \mathbf{1}_{\mathcal{N}_\lambda}$  converges to  $A(\mathbf{0}) \mathbf{1}_H$  in  $L^1(\mu)$ , and unravelling this statement we see that

$$\int_{\mathcal{N}_\lambda} A_\lambda(\mathbf{y}) e^{-S(\mathbf{y})} d\mathbf{y} \rightarrow \int_H A(\mathbf{0}) e^{-S(\mathbf{y})} d\mathbf{y}.$$

The last quantity is equal to  $c_0/2$ , showing that  $\lambda^{d/2} \mathcal{I}'(\lambda) \rightarrow c_0/2$  and finishing the proof.  $\square$

## 5. Proofs of main results

Theorem 2.3 differs from Theorem 4.1 in several ways. The most important is that the set where  $\Re \phi$  vanishes may extend to the boundary of the region of integration. This precludes the use of the easy deformation  $\pi$  because  $\mathcal{C}''$  is no longer supported on  $\{h > \epsilon\}$ . Consequently, some work is required to construct a suitable deformation. We do so via notions from stratified Morse theory [GM88].

**Tangent vector fields.** If  $\mathbf{x}$  is a point of the stratum  $S$  of the stratified space  $\mathcal{M}$ , let  $T_{\mathbf{x}}(\mathcal{M})$  denote the tangent space to  $S$  at  $\mathbf{x}$ . Because  $\mathcal{M}$  is embedded in  $\mathbb{C}^d$ , the tangent spaces may all be identified as subspaces of  $\mathbb{C}^d$ . Thus we have a notion of the tangent bundle  $T\mathcal{M}$ , a section of which is simply a vector field  $f$  on  $\mathcal{M} \subseteq \mathbb{C}^d$  such that  $f(\mathbf{x}) \in T_{\mathbf{x}}(\mathcal{M})$  for all  $\mathbf{x}$ . A consequence of the two Whitney conditions is the local product structure of a stratified space: a point  $\mathbf{p}$  in a  $k$ -dimensional stratum  $S$  of a stratified space  $\mathcal{M}$  has a neighborhood in which  $\mathcal{M}$  is homeomorphic to some product  $S \times X$ . According to [GM88], a proof may be found in mimeographed notes of Mather from 1970; it is based on Thom's Isotopy Lemma which takes up fifty pages of the same mimeographed notes. The next lemma is the only place where we use this (or any) consequence of Whitney stratification.

LEMMA 5.1. *Let  $f$  be a smooth section of the tangent bundle to  $S$ , that is  $f(s) \in T_s(S)$  for  $s \in S$ . Then each  $s \in S$  has a neighborhood in  $\mathcal{M}$  on which  $f$  may be extended to a smooth section of the tangent bundle.*

PROOF. Parametrize  $\mathcal{M}$  locally by  $S \times X$  and extend  $f$  by  $f(s, x) := f(s)$ .  $\square$

LEMMA 5.2 (vector field near a non-critical point). *Let  $\mathbf{x}$  be a point of the stratum  $S$  of the stratified space  $\mathcal{M}$  and suppose  $\mathbf{x}$  is not critical for the function  $\phi$ . Then there is a vector  $\mathbf{v} \in T_{\mathbf{x}}(S \otimes \mathcal{M})$  such that  $\operatorname{Re}\{d\phi(\mathbf{v})\} > 0$  at  $\mathbf{x}$ . Furthermore, there is a continuous section  $f$  of the tangent bundle in a neighborhood  $\mathcal{N}$  of  $\mathbf{x}$  such that  $\operatorname{Re}\{d\phi(f(\mathbf{y}))\} > 0$  at every  $\mathbf{y} \in \mathcal{N}$ .*

PROOF. By non-criticality of  $\mathbf{x}$ , there is a  $\mathbf{w} \in T_{\mathbf{x}}(S)$  with  $d\phi(\mathbf{w}) = u \neq 0$  at  $\mathbf{x}$ . Multiply  $\mathbf{w}$  componentwise by  $\bar{u}$  to obtain  $\mathbf{v}$  with  $\operatorname{Re}\{d\phi(\mathbf{v})\} > 0$  at  $\mathbf{x}$ . Use any chart map for  $S \otimes \mathbb{C}$  near  $\mathbf{x}$  to give a locally trivial coordinatization for the tangent bundle and define a section  $f$  to be the constant vector  $\mathbf{v}$ ; then  $\operatorname{Re}\{d\phi(f(\mathbf{y}))\} > 0$  on some sufficiently small neighborhood of  $\mathbf{x}$  in  $S$ . Finally, extend to a neighborhood of  $\mathbf{x}$  in  $\mathcal{M}$  by Lemma 5.1.  $\square$

Although we are working in the analytic category, the chains of integration are topological objects, for which we may use  $C^\infty$  methods (in what follows, even  $C^1$  methods will do). In particular, a partition of unity argument enhances the local result above to a global result.

LEMMA 5.3 (global vector field, in the absence of critical points). *Let  $\mathcal{M}$  be a compact stratified space and  $\phi$  a smooth function on  $\mathcal{M}$  with no critical points. Then there is a global section  $f$  of the tangent bundle of  $\mathcal{M}$  such that the real part of  $d\phi(f)$  is everywhere positive.*

PROOF. For each point  $\mathbf{x} \in \mathcal{M}$ , let  $f_{\mathbf{x}}$  be a section as in the conclusion of Lemma 5.2, on a neighborhood  $U_{\mathbf{x}}$ . Cover the compact space  $\mathcal{M}$  by finitely many sets  $\{U_{\mathbf{x}} : \mathbf{x} \in F\}$  and let  $\{\psi_{\mathbf{x}} : \mathbf{x} \in F\}$  be a smooth partition of unity subordinate to this finite cover. Define

$$f(\mathbf{y}) = \sum_{\mathbf{x} \in F} \psi_{\mathbf{x}}(\mathbf{y}) f_{\mathbf{x}}(\mathbf{y}).$$

Then  $f$  is smooth; it is a section of the tangent bundle because each tangent space is linearly closed; the real part of  $d\phi(f(\mathbf{y}))$  is positive because we took a convex combination in which each contribution was nonnegative and at least one was positive.  $\square$

Another partition argument gives our final version of this result.

**LEMMA 5.4** (global vector field, vanishing only at critical points). *Let  $\mathcal{M}$  be a compact stratified space and  $\phi$  a smooth function on  $\mathcal{M}$  with finitely many critical points. Then there is a global section  $f$  of the tangent bundle of  $\mathcal{M}$  such that the real part of  $d\phi(f)$  is nonnegative and vanishes only when  $\mathbf{y}$  is a critical point.*

**PROOF.** Let  $\mathcal{M}_\epsilon$  be the compact stratified space resulting in the removal of an  $\epsilon$ -ball around each critical point of  $\phi$ . Let  $f_\epsilon$  be a vector field as in the conclusion of Lemma 5.4 with  $\mathcal{M}$  replaced by  $\mathcal{M}_\epsilon$ . Let  $c_n$  be a positive real number, small enough so that the magnitudes of all partial derivatives of  $c_n f_{1/n}$  of order up to  $n$  are at most  $2^{-n}$ . In the topology of uniform convergence of derivatives of bounded order, the series  $\sum_n c_n f_n$  converges to a vector field  $f$  with the required properties.  $\square$

**Proof of Theorem 2.3.** Let  $f$  be a tangent vector field along which  $\phi$  increases away from critical points, as given by Lemma 5.4. Such a field gives rise to a differential flow, which, informally, is the solution to  $d\mathbf{p}/dt = f(\mathbf{p})$ . To be more formal, let  $\mathbf{x}$  be a point in a stratum  $S$  of  $\mathcal{M}$ . Via a chart map in a neighborhood of  $\mathbf{x}$ , we solve the ODE  $d\Phi(t)/dt = f(\Phi(t))$  with initial condition  $\Phi(0) = \mathbf{x}$ , obtaining a trajectory  $\Phi$  on some interval  $[0, \epsilon_{\mathbf{x}}]$  that is supported on  $S$ . Doing this simultaneously for all  $\mathbf{x} \in \mathcal{M}$  results in a map

$$\Phi : \mathcal{M} \times [0, \epsilon] \rightarrow \mathbb{C}^d$$

with  $\Phi(\mathbf{x}, t)$  remaining in  $S \otimes \mathbb{C}$  when  $\mathbf{x}$  is in the stratum  $S$ . The map  $\Phi$  satisfies  $\Phi(\mathbf{x}, 0) = \mathbf{x}$  and  $(d/dt)\Phi(\mathbf{x}, t) = f(\Phi(\mathbf{x}, t))$ . The fact that this may be defined up to time  $\epsilon$  for some  $\epsilon > 0$  is a consequence of the fact that the vector field  $f$  is bounded and that a small neighborhood of  $\mathcal{M}$  in  $\mathcal{M} \otimes \mathbb{C}$  is embedded in  $\mathbb{C}^d$ . Because  $f$  is smooth and bounded, for sufficiently small  $\epsilon$  the map  $\mathbf{x} \mapsto \Phi(\mathbf{x}, \epsilon)$  is a diffeomorphism.

The flow reduces the real part of  $\phi$  everywhere except the critical points which are rest points. Consequently, it defines a homotopy  $H(\mathbf{x}, t) := \Phi(\mathbf{x}, \epsilon t)$  between  $\mathcal{C}$  and a chain  $\mathcal{C}'$  on which the minima of the real part of  $\phi$  occur precisely on the set  $G$ . Recall that  $H$  induces a chain homotopy  $\mathcal{C}_H$  with  $\partial\mathcal{C}_H = \mathcal{C}' - \mathcal{C} + \partial\mathcal{C} \times \sigma$ , where  $\sigma$  is a standard 1-simplex. Let  $\omega$  denote the holomorphic  $d$ -form  $A(\mathbf{z}) \exp(-\lambda\phi(\mathbf{z})) d\mathbf{z}$ . Because  $\omega$  is a holomorphic  $d$ -form in  $\mathbb{C}^d$ , we have  $d\omega = 0$ . Now, by Stokes' Theorem,

$$\begin{aligned} 0 &= \int_{\mathcal{C}_H} d\omega \\ &= \int_{\partial\mathcal{C}_H} \omega \\ &= \int_{\mathcal{C}'} \omega - \int_{\mathcal{C}} \omega - \int_{\partial\mathcal{C} \times \sigma} \omega. \end{aligned}$$

The chain  $\partial\mathcal{C} \times \sigma$  is supported on a finite union of spaces  $S \otimes \mathcal{C}$  where  $S$  is a stratum of dimension at most  $d - 1$ . The integral of a holomorphic  $d$ -form vanishes over such a chain. Therefore, the last term on the right drops out and we have

$$\int_{\mathcal{C}} \omega = \int_{\mathcal{C}'} \omega.$$

Outside of a neighborhood of  $G$  the magnitude of the integrand is exponentially small, so we have shown that there are  $d$ -chains  $\mathcal{C}_{\mathbf{x}}$  supported on arbitrarily small

neighborhoods  $\mathcal{N}(\mathbf{x})$  of each  $\mathbf{x} \in G$  such that

$$(5.1) \quad \mathcal{I}(\lambda) - \sum_{\mathbf{x} \in G} \int_{\mathcal{C}_{\mathbf{x}}} \omega$$

is exponentially small. To finish that proof, we need only show that each  $\int_{\mathcal{C}_{\mathbf{x}}} \omega$  has an asymptotic series in decreasing powers of  $\lambda$  whose leading term, when  $A(\mathbf{x}) \neq 0$ , is given by

$$(5.2) \quad c_0(\mathbf{x}) = (2\pi)^{-d/2} A(\mathbf{x}) e^{\lambda \phi(\mathbf{x})} (\det \mathcal{H}(\mathbf{x}))^{-1/2}.$$

The  $d$ -chain  $\mathcal{C}_{\mathbf{x}}$  may be parametrized by a map  $\psi_{\mathbf{x}} : B \rightarrow \mathcal{N}(\mathbf{x})$ , mapping the origin to  $\mathbf{x}$ , where  $B$  is the open unit ball in  $\mathbb{R}^d$ . By the chain rule,

$$\int_{\mathcal{C}_{\mathbf{x}}} \omega = \int_B [A \circ \psi](\mathbf{x}) \exp(-\lambda[\phi \circ \psi(\mathbf{x})]) \det d\psi(\mathbf{x}) d\mathbf{x}.$$

The real part of the analytic phase function  $\phi \circ \psi$  has a strict minimum at the origin, so we may apply Theorem 4.1. We obtain an asymptotic expansion whose first term is

$$(5.3) \quad (2\pi\lambda)^{-d/2} [A \circ \psi](\mathbf{0}) \det d\psi(\mathbf{0}) (\det M_{\mathbf{x}})^{-1/2}$$

where  $M_{\mathbf{x}}$  is the Hessian matrix of the function  $\phi \circ \psi$ . The term  $[A \circ \psi](\mathbf{0})$  is equal to  $A(\mathbf{x})$ . The Hessian matrix of  $\phi \circ \psi$  at the origin is given by  $M_{\mathbf{x}} = d\psi(\mathbf{0}) \mathcal{H}(\mathbf{x}) d\psi(\mathbf{0})$ . Thus

$$\det M_{\mathbf{x}} = (\det d\psi(\mathbf{0}))^2 \det \mathcal{H}(\mathbf{x})$$

and plugging into (5.3) yields (5.2).  $\square$

Proof of Corollary 2.4: Lemma 5.4 does not require the critical points to be in the interior, so the argument leading up to (5.1) is still valid. For those points  $\mathbf{x}$  in a  $(d-1)$ -dimensional stratum, use Corollary 2.4 in place of Theorem 4.1 to obtain (5.2) with an extra factor of  $1/2$ .  $\square$

REMARK. *The reason we do not continue with a litany of special geometries (quarter-spaces, octants, and so forth) is that the case of a halfspace is somewhat special. The differential of the change of variables at the origin is a nonsingular map, which must send half-spaces to half-spaces, though it will in general alter angles of any smaller cone.*

## 6. Examples

The simplest multidimensional application of our results is a computation from [PW04]. The purpose is to estimate coefficients of a class of bivariate generating functions whose denominator is the product of two smooth divisors. We give only a brief summary of how one arrives at (6.1) from a problem involving generating functions; a complete explanation of this can be found in [PW04, Section 4]. Note, however, that the mathematics of the integral is not contained in that paper, which instead refers to an earlier draft of this one!

Let  $v_1, v_2$  be distinct analytic functions of  $z$  with  $v_1(1) = v_2(1) = 1, 0 \neq v'_1(1) \neq v'_2(1) \neq 0$  and such that each  $|v_i(z)|$  attains its maximum on  $|z| = 1$  only at  $z = 1$ . For example, the last condition is satisfied by any pair of aperiodic power series with nonnegative coefficients and radius of convergence greater than 1.

Consider the generating function  $F(z, w) = 1/H(z, w)$ , where  $H(z, w) = (1 - wv_1(z))(1 - wv_2(z))$ . The two branches of the curve  $H = 0$  intersect only at  $(1, 1)$  and this intersection is transverse. The Maclaurin coefficients of  $F(z, w) = \sum_{r,s} a_{rs} z^r w^s$  are given by the Cauchy integral formula

$$a_{rs} = \frac{1}{(2\pi i)^2} \frac{dw dz}{z^{r+1} w^{s+1} (1 - wv_1(z))(1 - wv_2(z))}$$

where the integral is taken over a product of circles centred at  $(0, 0)$  and of sufficiently small radii.

Pushing the contour out to  $|z| = 1, |w| = 1 - \varepsilon$  we obtain the same formula, since  $F$  is still analytic inside the product of disks bounded by these latter circles. Pushing the  $w$ -contour out to  $|w| = 1 + \varepsilon$ , using the residue formula on the inner integral and observing that the integral over  $|w| = 1 + \varepsilon$  is exponentially decaying as  $s \rightarrow \infty$ , we see that

$$a_{rs} \approx \frac{1}{2\pi} \int_{|z|=1} \frac{-R_s(z)}{z^{r+1}} dz$$

where  $\approx$  means that the difference is exponentially decaying as  $s \rightarrow \infty$  and  $R_s(z)$  denotes the sum of residues of  $w \mapsto w^{-(s+1)} F(z, w)$  at the roots  $w = 1/v_i(z)$ ,  $i \in \{1, 2\}$ .

The residue sum  $R_s(z)$  can be rewritten in terms of an integral via

$$-R_s(z) = (s+1) \int_0^1 [(1-p)v_1 + pv_2]^s dp$$

and so we have

$$a_{rs} \approx \frac{s+1}{2\pi} \int_{|z|=1} z^{-(r+1)} \int_0^1 [(1-p)v_1(z) + pv_2(z)]^s dp dz.$$

In order to cast this into our standard framework, we need to be able to define a branch of the logarithm of  $(1-p)v_1(z) + pv_2(z)$ . We do this by localizing on the circle  $|z| = 1$  to a sufficiently small neighbourhood of the point  $z = 1$ . This is possible since the integrand decays exponentially away from  $z = 1$ , by hypotheses on the  $v_i$ , and we shall show that the integral near  $z = 1$  decays only polynomially.

The substitution  $z = e^{i\theta}$  converts this to an integral

$$(6.1) \quad a_{rs} \approx \frac{s+1}{2\pi} \int_{\mathcal{N}} \int_0^1 e^{-s\phi(p,t)} A(p, z) dp dt$$

where  $\phi(p, t) = ir\theta/s + \log [(1-p)v_1(e^{it}) + pv_2(e^{it})]$ ,  $A(p, t) = 1$ , and  $\mathcal{N}$  is a closed interval centred at 0. To compute asymptotics in the direction  $r/s = \kappa$ , for fixed  $\kappa > 0$ , we can consider  $\phi$  to be independent of  $r$  and  $s$ .

We now asymptotically evaluate (6.1) using Theorem 2.3. We can rewrite the iterated integral as a single integral over the stratified space  $\mathcal{M} = \mathcal{N} \times [0, 1]$ . The phase  $\phi$  has nonnegative real part and this fits into our framework. There is a single stationary point, at  $(p, z) = (1/2, 0)$  (note that  $\text{Re}\{\phi\}$  is zero for all  $(p, 0)$ , so Theorem 4.1 does not suffice). This critical point is quadratically nondegenerate and direct computation using Theorem 2.3 yields

$$(6.2) \quad a_{rs} = \frac{1}{|v'_1(1) - v'_2(1)|} + O(s^{-1})$$

as  $s \rightarrow \infty$  with  $\kappa$  fixed. By keeping track of error terms more explicitly, it is easily shown that this approximation is uniform in  $\kappa$  provided  $\kappa$  stays in a compact subset

of the open interval formed by  $v'_1(1), v'_2(1)$  (it follows from our assumptions that these numbers are positive real — see [PW04] for more details). This means that  $a_{rs}$  is asymptotically constant in any compact subcone of directions away from the boundary formed by the lines  $\kappa_i = v'_i(1)$ .

This example, and in fact a number of cases in [PW04], can also be solved using iterated residues. This is carried out in [BP04]. Iterated residues have the advantage of showing that the  $O(s^{-1})$  term decays exponentially, but the disadvantage that they do not give any results when  $\kappa$  approaches the boundary. The present methods do give boundary results. Corollary 2.4 shows that  $a_{rs}$  converges to one half the right-hand side of (6.2) when  $(r, s) \rightarrow \infty$  with  $r/s = \kappa_1 + O(1)$ , and a small extension obtains a Gaussian limit: letting  $\Phi$  denote the standard normal cumulative distribution function, we have

$$a_{rs} = \frac{\Phi(u)}{|v'_1(1) - v'_2(1)|} + O(s^{-1})$$

when  $r, s \rightarrow \infty$  with  $(r/s - \kappa_1)/s^{1/2} \rightarrow u$ .

## 7. Further comments

We have not emphasized explicit formulae for the higher order terms, giving an equation such as (2.1) only for the leading term in the case where  $A(\mathbf{0}) \neq 0$ . However, our results establish the validity of existing computations of higher order terms under our more general hypotheses.

To elaborate, we prove Theorem 2.3 by first constructing a change of variables  $\mathbf{x} \mapsto \Phi(x, \epsilon)$  homotopic to the identity under which the minimum of  $\text{Re}\{\phi\}$  at  $\mathbf{0}$  is strict, and then changing variables, again homotopically to the identity, to the standard form. The composition  $\psi$  of these two maps is homotopic to the identity but is far from explicitly given: while the second map is constructed by an explicitly defined Morse function, the first deformation is the solution to a differential equation and is not particularly explicit.

In [Hör90], Hörmander derives such an explicit formula (assuming smoothness) for integrals of our type where  $\mathcal{M} = \mathbb{R}^d$  and  $A$  has compact support. The formula is indeed rewritten and used in [RW08] to compute higher order terms for generating function applications, in which more restrictive hypotheses preclude the vanishing of  $\text{Re}\{\phi\}$  on a curve reaching the boundary of the chain of integration. Their methods, while not covering the cases of interest here, do have the virtue of dealing with the change of variables  $\psi$  only through the equation  $S = \phi \circ \psi$ . In particular, the derivatives of  $\psi$  arising in the computation of the new amplitude function  $(A \circ \psi) \det d\psi$  can be computed by implicitly differentiating the equation  $S = \phi \circ \psi$ . Having found at least one such  $\psi$  homotopic to the identity, we are now free to replicate the computations of [RW08] under our more general hypotheses, as follows.

In the case of standard phase, the coefficient of  $\lambda^{-(n+d)/2}$  is given by (provided all  $r_i$  are even)

$$\sum_{|\mathbf{r}|=n} a_{\mathbf{r}} \beta_{\mathbf{r}}$$

where  $a_{\mathbf{r}}$  is the Maclaurin coefficient of  $A$  corresponding to the monomial  $\mathbf{r}$  and  $\beta_{\mathbf{r}}$  is the constant defined in Corollary 3.2. Note that  $n$  must be even for this coefficient to be nonzero, so we write  $n = 2k$ . The differential operator  $\partial_1^{r_1} \dots \partial_d^{r_d}$



when applied to  $A$  and evaluated at  $\mathbf{0}$  yields precisely  $\prod_i r_i! a_{\mathbf{r}}$ . Thus the operator

$$\sum_{|\mathbf{r}|=k} \frac{\partial_1^{2r_1} \dots \partial_d^{2r_d}}{4^k r_1! \dots r_d!}$$

applied to  $A$  and evaluated at zero yields the coefficient we seek.

After the Morse lemma is applied using the change of variables  $S = \phi \circ \psi$ , we need to apply the displayed operator to the new amplitude  $(A \circ \psi) \det d\psi$ . The resulting expression evaluated at  $\mathbf{x}$  can be computed directly via the rules of Leibniz and Faa di Bruno. Evaluating at  $\mathbf{x}$  simplifies some terms, and, as mentioned above, derivatives of  $(A \circ \psi) \det d\psi$  may be computed without explicitly specifying  $\psi$ .

As a relatively simple example, consider the case  $k = 1$  and  $d = 1$ . The differential operator reduces to  $\frac{1}{4}\partial^2$  where  $\partial$  denotes differentiation with respect to the variable  $x$ . Applying this to  $(A \circ \psi) \det d\psi$  yields (with superscripts denoting the order of derivatives and arguments suppressed)

$$\frac{1}{4} \left( A^{(2)} \left( \psi^{(1)} \right)^3 + 3A^{(1)} \psi^{(1)} \psi^{(2)} + A^{(0)} \psi^{(3)} \right).$$

The defining equation  $S = \phi \circ \psi$  can be differentiated to yield the system

$$\begin{aligned} 2x &= \phi^{(1)} \psi^{(1)} \\ 2 &= \phi^{(2)} \left[ \psi^{(1)} \right]^2 + \phi^{(1)} \psi^{(2)} \\ 0 &= \phi^{(3)} \left[ \psi^{(1)} \right]^3 + 3\phi^{(2)} \psi^{(1)} \psi^{(2)} + \phi^{(1)} \psi^{(3)} \\ 0 &= \phi^{(4)} \left[ \psi^{(1)} \right]^4 + 6\phi^{(3)} \left[ \psi^{(1)} \right]^2 \psi^{(2)} + 4\phi^{(2)} \psi^{(1)} \psi^{(3)} + 3\phi^{(2)} \left[ \psi^{(2)} \right]^2 + \phi^{(1)} \psi^{(4)}. \end{aligned}$$

Evaluating these at the point in question, we see that the terms with highest derivatives of  $\psi$  vanish in each equation. The system is triangular and can be solved explicitly to obtain

$$\begin{aligned} \psi^{(1)} &= \sqrt{\frac{2}{\phi^{(2)}}} \\ \psi^{(2)} &= \frac{-2\phi^{(3)}}{3 \left[ \phi^{(2)} \right]^2} \\ \psi^{(3)} &= \frac{\left[ 5\phi^{(3)} \right]^2 - 3\phi^{(2)} \phi^{(4)}}{3\sqrt{2} \left[ \phi^{(2)} \right]^{7/2}}. \end{aligned}$$

Putting these together with the expression for the derivative of  $(A \circ \psi) \det d\psi$  above yields an expression for the  $\lambda^{-3/2}$  term in the integral that is a rational function with denominator  $\left[ \phi^{(2)} \right]^{7/2}$  and numerator a polynomial in the derivatives of  $A$  up to order 2 and  $\phi$  to order 4. In summary, the results of this paper show that the computational apparatus and formulae for higher order terms given in [RW08] hold in the case of complex phase functions integrated over stratified spaces.

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